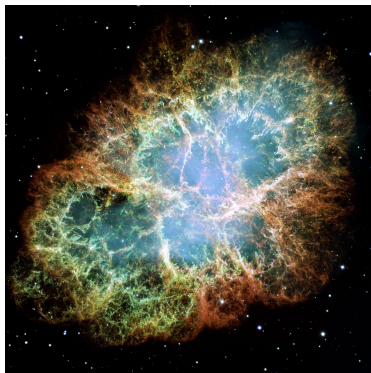
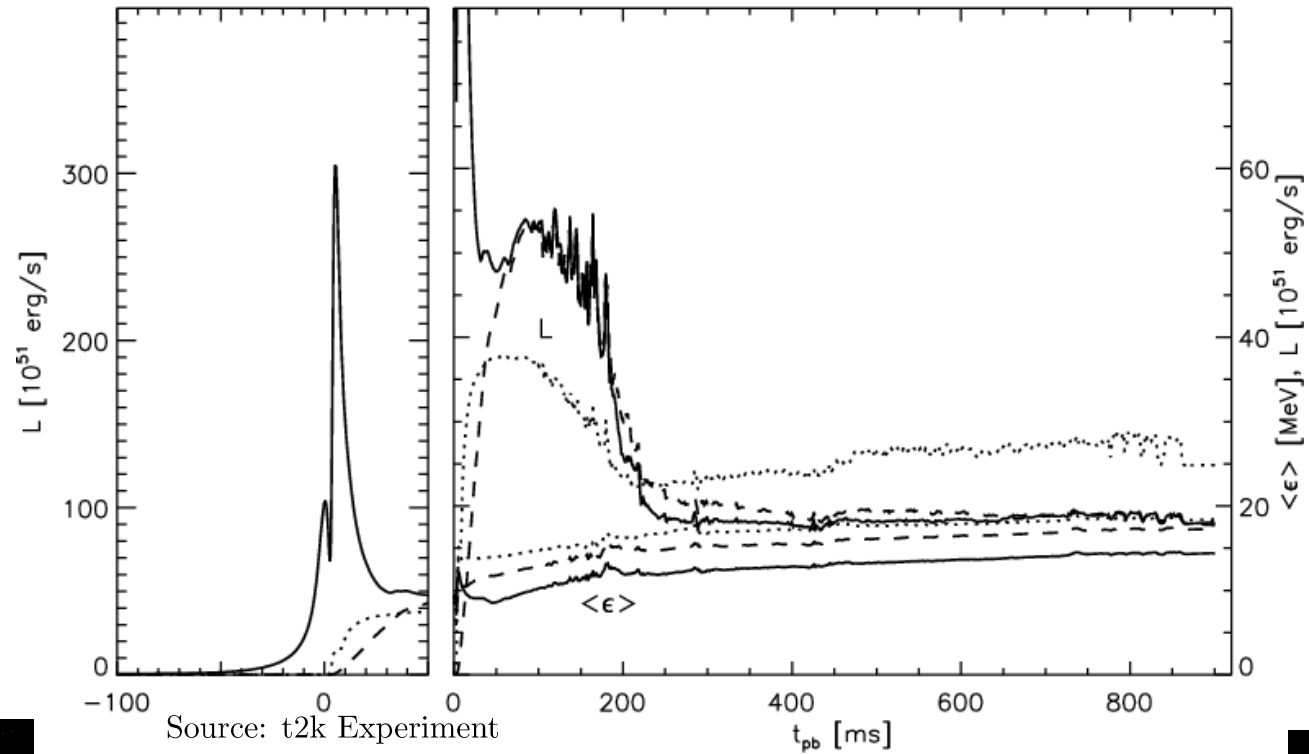


Neutrino Transport in Core-Collapse Supernovae

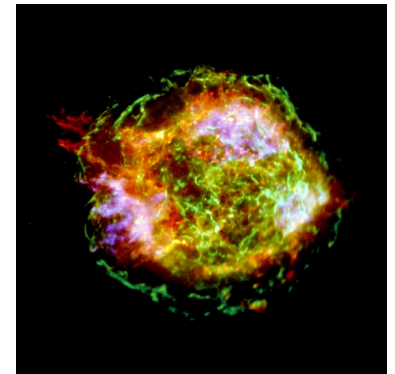


Crab Nebula (Hubble Space Telescope)

Master's Presentation Talk
Paul Mekhedjian
October 23rd, 2014



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DARMSTADT



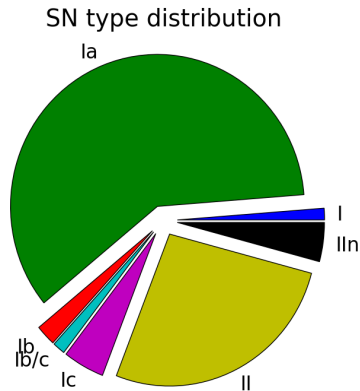
Cassiopeia A (Chandra X-Ray Observatory)

Talk Outline

- Introduction & Background
- Goals & Methods
- Results & Analysis
- Summary & Conclusion

Core Collapse Death of Massive Stars

- Focus on Type II SNe:



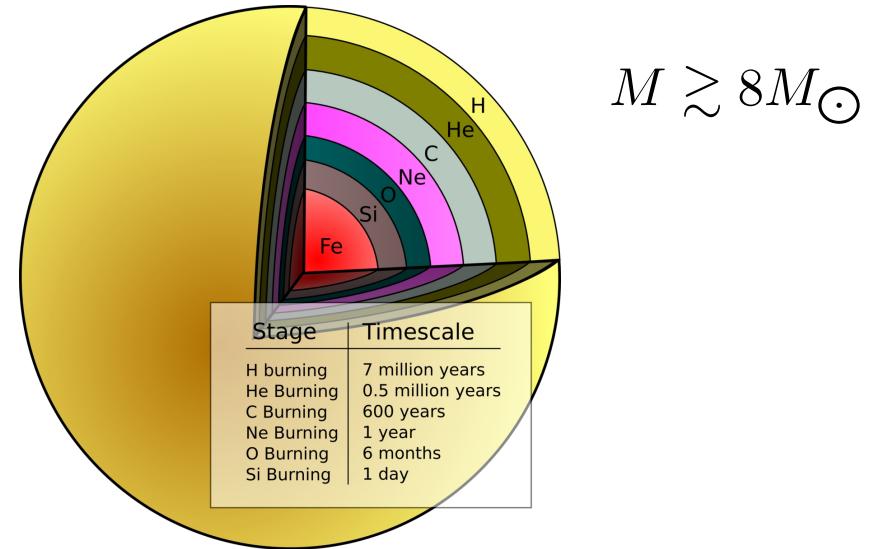
Source: Padova-Asiago Supernova Group

Type Ia: Thermonuclear explosion
Types II, Ib/Ic: Core Collapse

- Chandrasekhar mass defines maximum stable mass supported by degenerate electron pressure:

$$M_{\text{ch}} \simeq 1.457 \cdot \left(\frac{Y_e}{0.5} \right)^2 \cdot M_{\odot}$$

- Burning at exponential timescales...



Source: A. C. Phillips, *The Physics of Stars*

- Iron core conditions right before / at collapse:

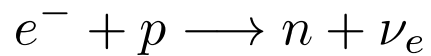
$$\rho_c \simeq 10^9 - 10^{10} \text{ g/cm}^3$$

$$T_c \simeq (8 - 10) \times 10^9 \text{ K}$$

$$t_{\text{collapse}} \simeq \frac{1}{\sqrt{G \cdot \rho_c}} \longrightarrow t_{\text{collapse}} \approx 100 \text{ ms}$$

Collapse Dynamics

1. As ρ_c increases, so does μ_e : $\mu_e(\rho) \propto \rho^{1/3}$



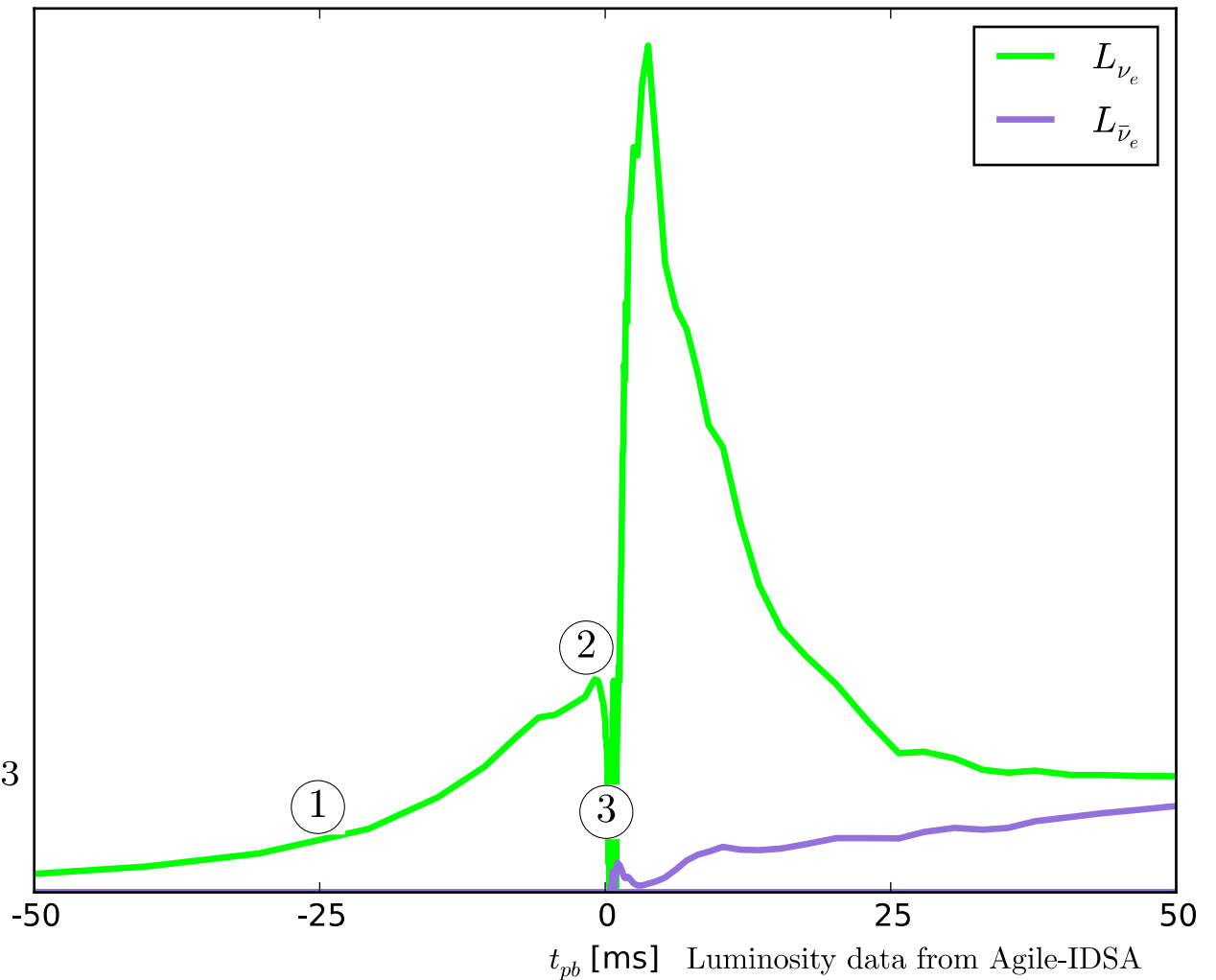
Electron capture on bound protons lowers Y_e and consequently M_{ch} and collapse of iron core ensues.

2. At $\rho \gtrsim 2 \times 10^{12} \text{ g/cm}^3$ neutrinos begin to become trapped:

$$\tau_{\text{diffusion}} \gtrsim \tau_{\text{collapse}}$$

3. Nuclear saturation density at $\rho \simeq 2.7 \times 10^{14} \text{ g/cm}^3$ leads to stronger neutrino diffusion, core bounce and shock:

$$E_{\text{shock}}^{\text{init}} \approx (4 - 10) \times 10^{51} \text{ erg}$$



Neutrino Burst, Accretion and Heating

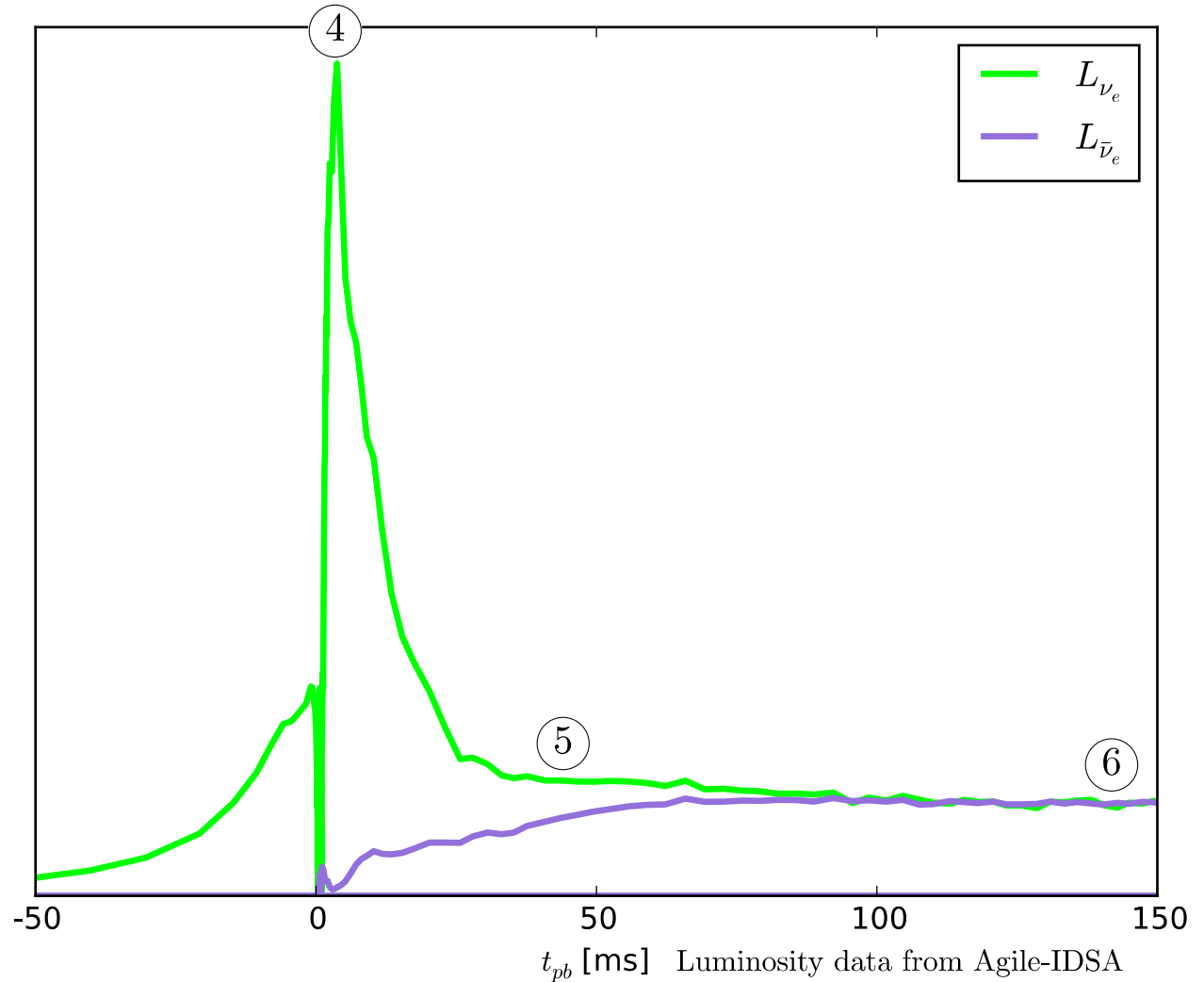
4. Neutrino burst at ~ 3 ms post-bounce. Deleptonization and ν -burst weaken shock. Photo-dissociation of baryonic matter eventually stops it (*prompt shock stalling*).



5. Gravitational energy from accretion fuels added neutrino luminosity: L_{ν}^{acc}

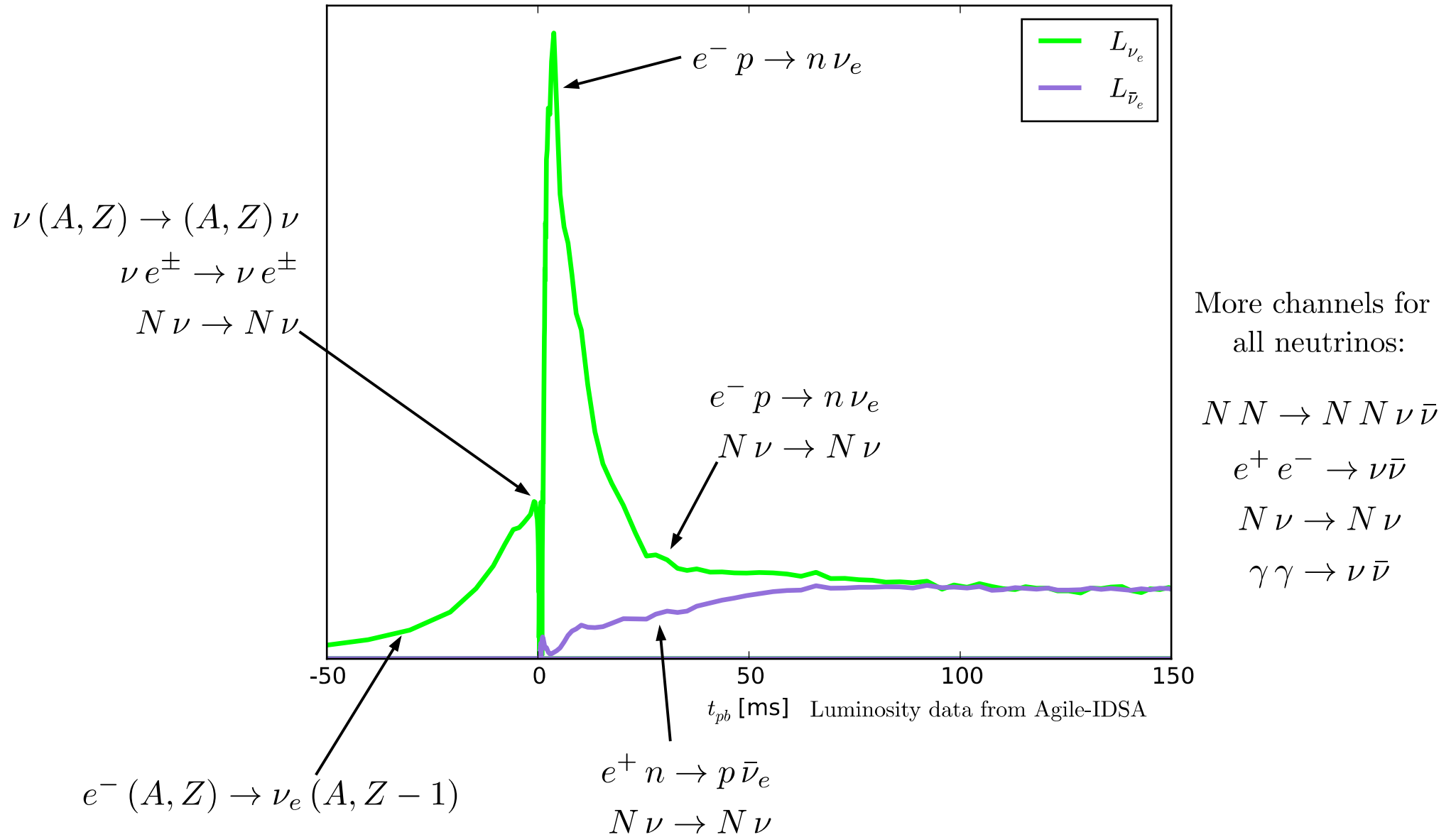
$T \uparrow, \rho \downarrow$ yields $\bar{\nu}_e, \nu_{\mu/\bar{\mu}}, \nu_{\tau/\bar{\tau}}$

6. Neutrino absorption and continuation of neutrino-driven mechanism (*delayed shock revival*):



- Neutrino shock revival mechanism: Wilson(1985), Bethe (1990)
~ 1% - 10% of neutrinos are enough to revive shock

ν -relevant Reactions



Goal & General Strategy

Goal:

- Develop a new neutrino treatment which is simple, efficient, and reproduces most important features of neutrinos in core collapse supernova (CCSN) environment

Strategy:

- ▶ Explore CCSN dynamics with Agile-IDSA (1D, spherically symmetric GR code) using $15M_{\odot}$ progenitor, Lattimer-Swesty nuclear EOS
- ▶ Use 1D profiles (ρ, T, Y_e) , develop our scheme to obtain $\dot{\epsilon}$ and \dot{Y}_e
- ▶ Test scheme against known quantities such as L_{ν} , $\langle \epsilon_{\nu} \rangle$, $\dot{\epsilon}$, \dot{Y}_e for reproducibility

Code Integration:

- ▶ Replacement as a candidate for neutrino transport in hydrodynamics codes via $\dot{\epsilon}$ and \dot{Y}_e
- ▶ Ultimately, be able to use scheme for long-time simulations of neutron star cooling and nucleosynthesis

Building Energy & Abundance Rates from Radial Derivatives of Luminosity

- After bounce, for $r > r_{\text{trapping}} \equiv R_\nu$:

$$\dot{e} = - \left[\frac{\partial L_{\nu_e}}{\partial r} + \frac{\partial L_{\bar{\nu}_e}}{\partial r} \right] \cdot \frac{1}{4\pi r^2 \rho}$$

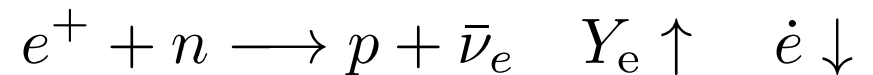
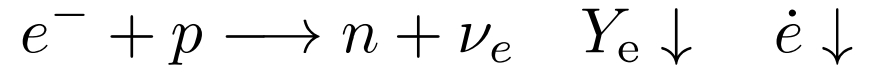
$$\dot{Y}_e = \left[\frac{\partial L_{\bar{\nu}_e}^N}{\partial r} - \frac{\partial L_{\nu_e}^N}{\partial r} \right] \cdot \frac{m_b}{4\pi r^2 \rho}$$

- Construct infinitesimal luminosity from heating and cooling, based on Janka (2001):

$$\frac{1}{4\pi r^2 \rho} \frac{\partial L_\nu(r)}{\partial r} = - q_\nu^+(r) + q_\nu^-(r)$$

$$\frac{1}{4\pi r^2 \rho} \frac{\partial L_\nu^N(r)}{\partial r} = - r_\nu^+(r) + r_\nu^-(r)$$

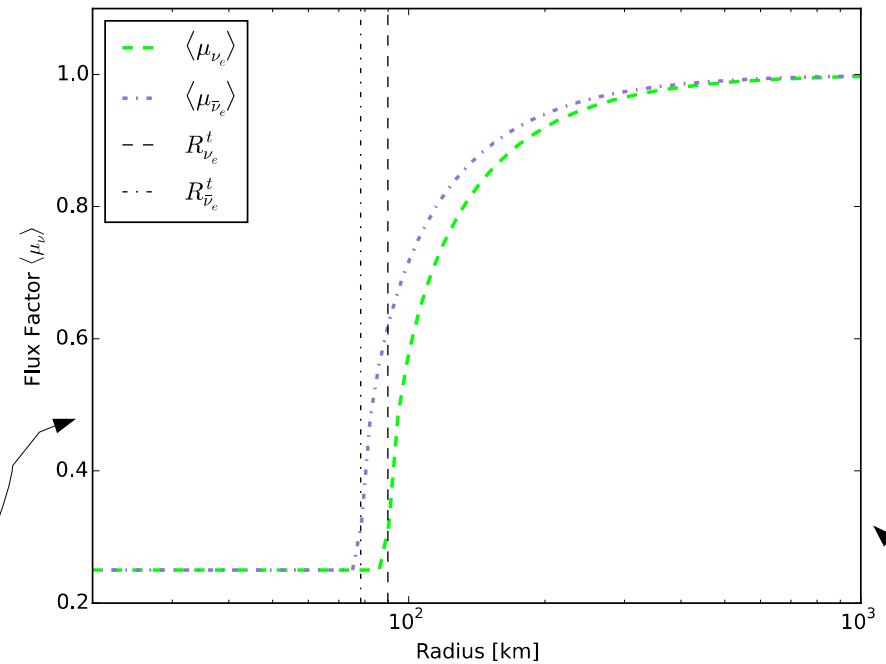
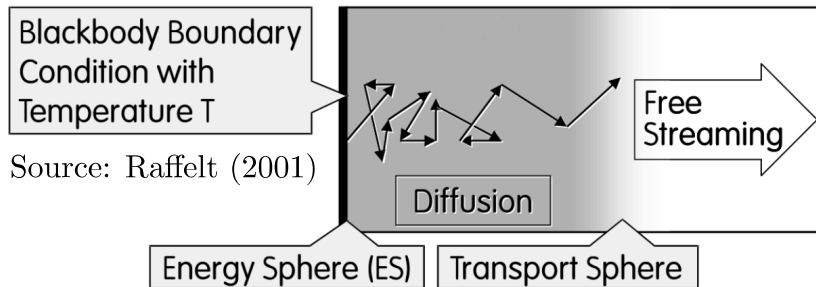
- Dominant reactions for ν_e & $\bar{\nu}_e$:



where q_ν is the specific energy rate in
[erg/g · s] and r_ν is the specific number
rate in [# / g · s]

Luminosity Estimation Scheme

New Ansatz - Based on Janka (2001):



$$L_\nu^N(r) = L_\nu^N(R_\nu) \cdot \exp \left\{ - \int_{R_\nu}^r dr' \frac{\langle \kappa_\nu^{\text{abs}} \rangle}{\langle \mu_\nu \rangle} \right\} + \int_{R_\nu}^r dr' \cdot 4\pi(r')^2 \cdot R_{\nu,\text{cool}}^- \cdot \exp \left\{ - \int_{r'}^r dr'' \frac{\langle \kappa_{\text{abs}} \rangle}{\langle \mu_\nu \rangle} \right\}$$

dampening
re-absorption

- Energy luminosity determined by:

$$L_\nu(r) \simeq \langle \epsilon_\nu \rangle \Big|_{r=R_\nu} \cdot L_\nu^N(r)$$

$$\langle \epsilon_\nu \rangle \Big|_{r=R_\nu} = T(R_\nu) \cdot \overbrace{\frac{\mathcal{F}_3(0)}{\mathcal{F}_2(0)}}^{\sim 3.15\dots} \simeq 3.15\dots \cdot T(R_\nu)$$

Neutrino spheres

Motivation: Identify the location of the boundary condition

- Cross Section

$$\sigma_\nu \propto \sigma_0 \cdot \epsilon_\nu^2 \rightarrow \kappa_\nu = \frac{\sigma_\nu \cdot \rho \cdot Y_i}{m_b}$$

- Total Opacity

$$\kappa_\nu^t = \kappa_\nu^{\text{abs}} + \kappa_\nu^{\text{scat}}$$

- Effective Opacity

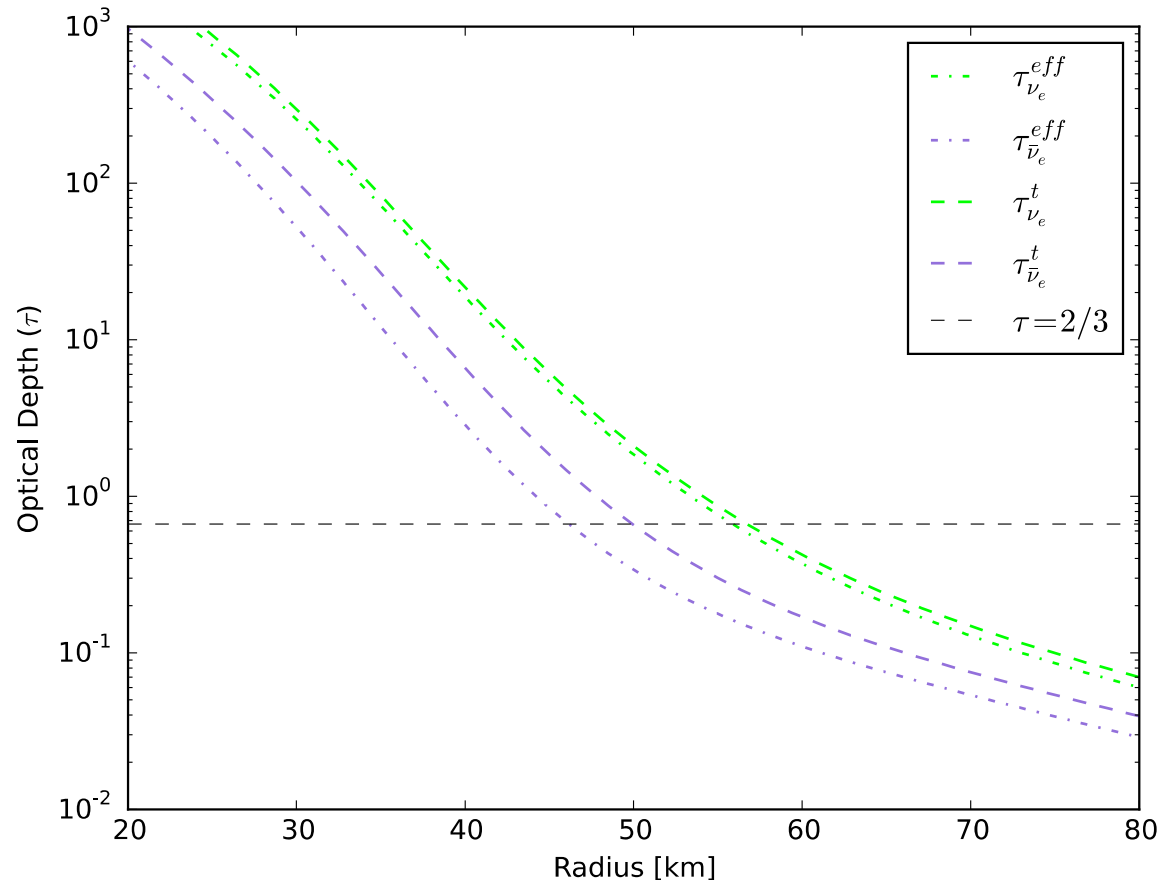
$$\kappa_\nu^{\text{eff}} = \sqrt{\kappa_\nu^{\text{abs}} \cdot (\kappa_\nu^{\text{abs}} + \kappa_\nu^{\text{scat}})}$$

- Optical Depth & Neutrino spheres

$$\tau_\nu(r) = \int_r^\infty \kappa_\nu(r') dr'$$

$$\tau_\nu(R_\nu) = 2/3 = \int_{R_\nu}^\infty \kappa_\nu(r') dr'$$

- τ_ν may be physically interpreted as number of ν interactions before leaving system (i.e. neutrino at $\tau_\nu \simeq 100$ will experience ~ 100 interactions on average of given type)



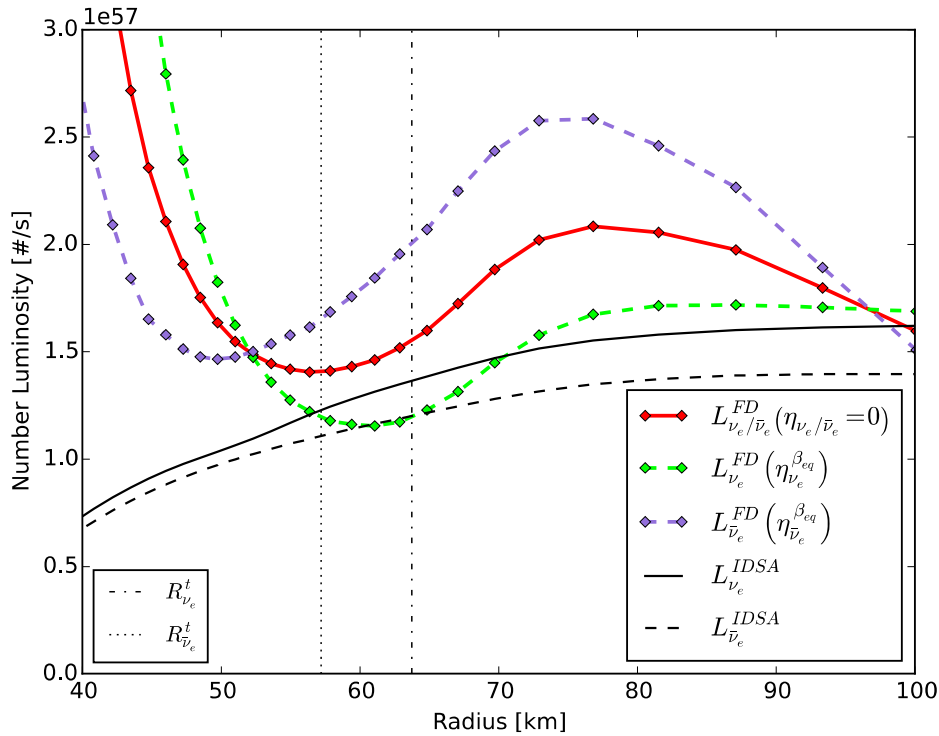
Fermi Blackbody Luminosity

Motivation: Characterize nature of and give a value to the boundary condition

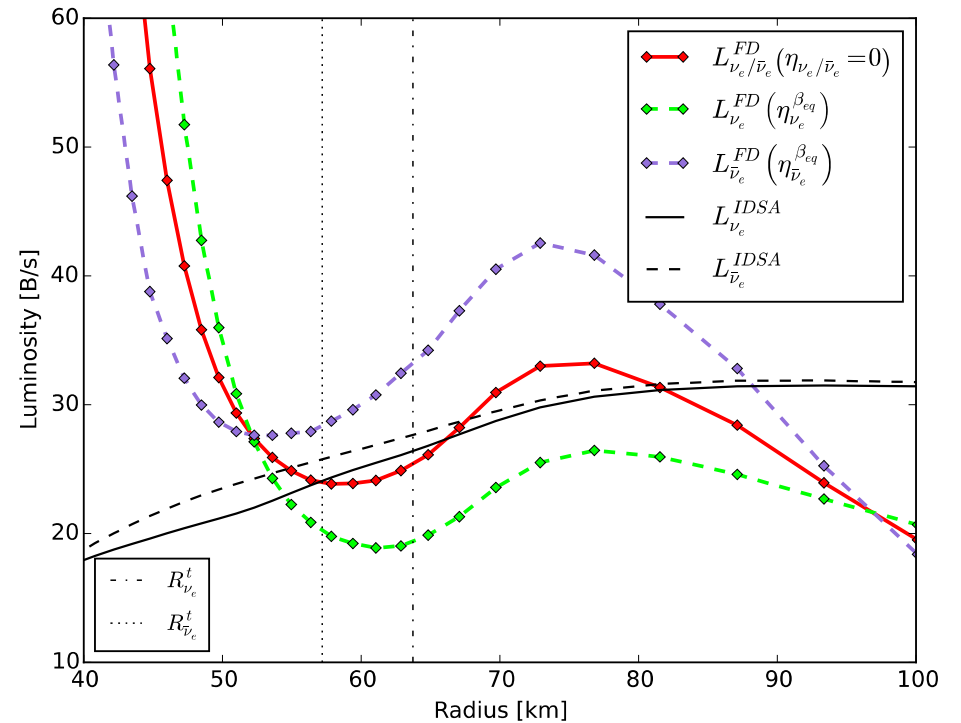
- Mimic production of neutrinos via blackbody luminosity. Near R_ν , $T \simeq T_\nu$.
- Test ν degeneracy at neutrino spheres. Which hierarchy is best assumption?

$$L_\nu^{N,FD} = \frac{4\pi^2 \cdot c}{(hc)^3} \cdot R_\nu^2 \cdot (k_B T_\nu)^3 \cdot \mathcal{F}_2(\eta_\nu)$$

$$L_\nu^{FD} = \frac{4\pi^2 \cdot c}{(hc)^3} \cdot R_\nu^2 \cdot (k_B T_\nu)^4 \cdot \mathcal{F}_3(\eta_\nu)$$



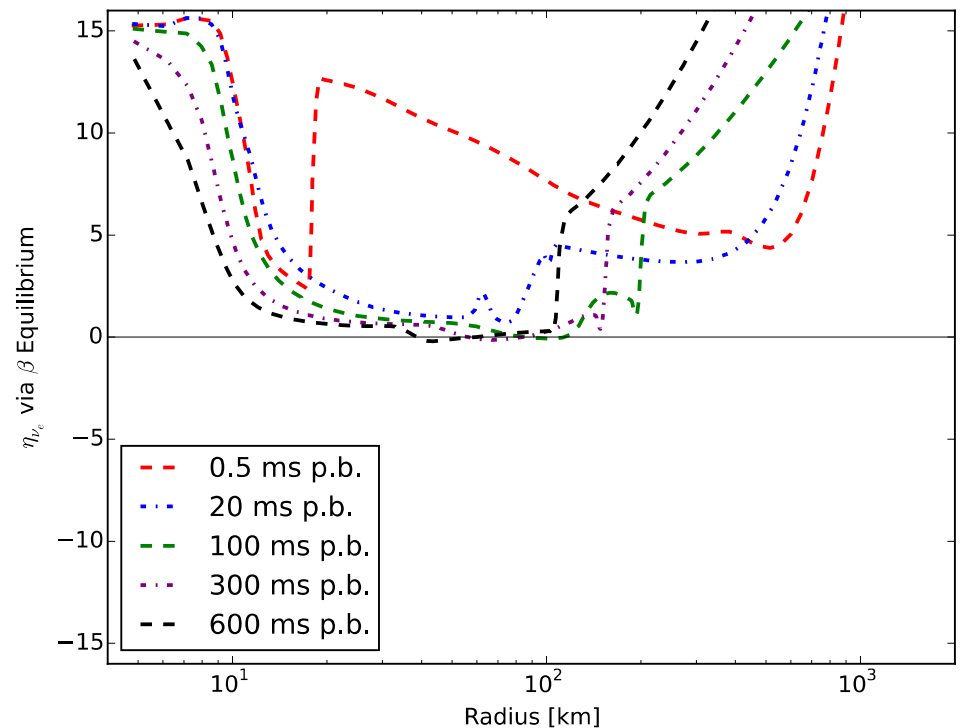
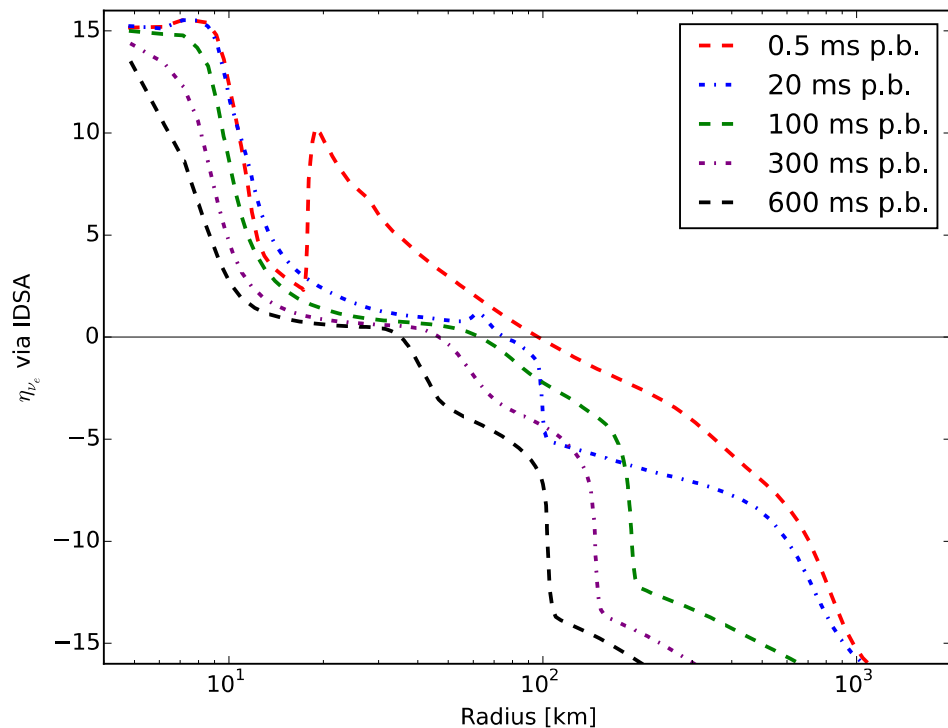
$t_{\text{pb}} \sim 300$ ms p.b.



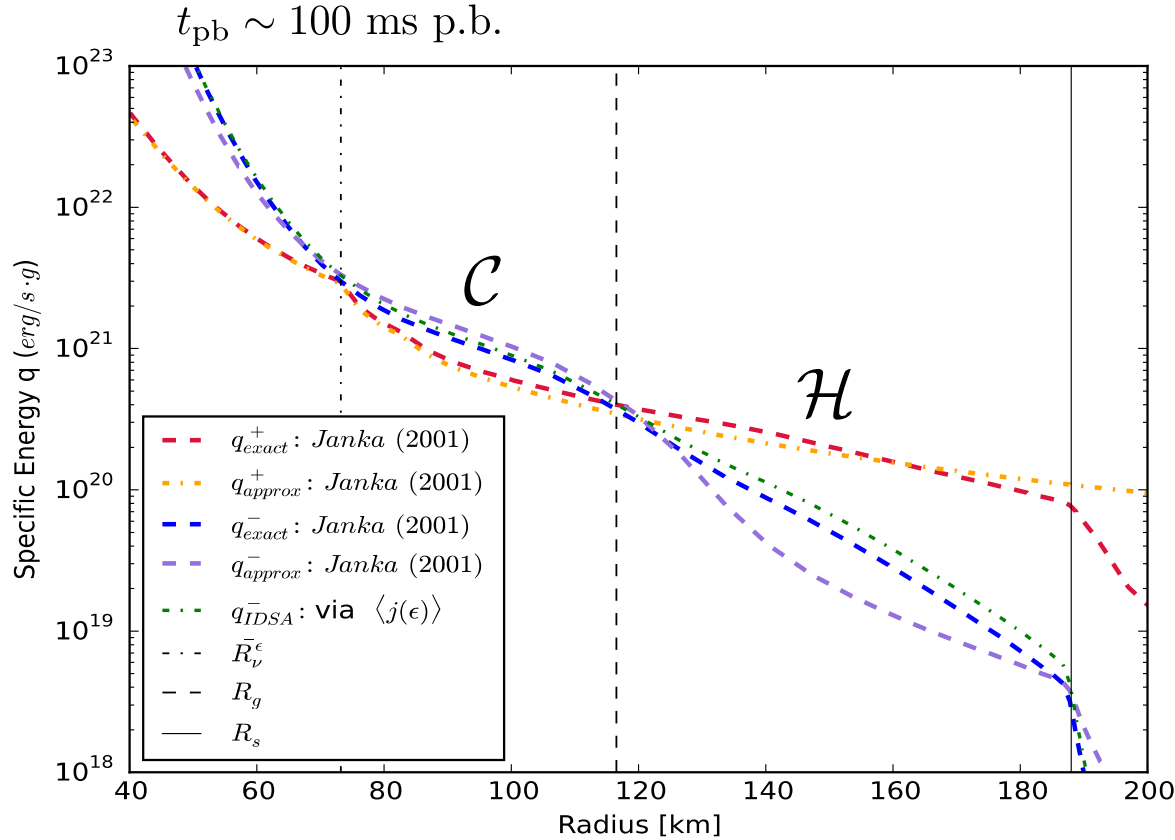
$t_{\text{pb}} \sim 300$ ms p.b.

Neutrino Degeneracy at Various Times

- Comparison shows excellent agreement in β equilibrium: $\rho \gtrsim 10^{11} - 10^{12} \text{ g/cm}^3$
- $e^- + p \longrightarrow n + \nu_e$ yields $\mu_e + \mu_p \longrightarrow \mu_n + \mu_{\nu_e}$
 $\mu_{\nu_e} \longrightarrow +\mu_e - \hat{\mu} - \Delta_{np}, \quad \mu_{\bar{\nu}_e} = -\mu_{\nu_e}$
- Assumption of $\eta_\nu = 0$ everywhere exclusively is insufficient for neutrino transport, but seems to be enough for Fermi blackbody boundary condition at later times! Used by Janka (2001) in: $R_\nu < r < R_s$



Neutrino Heating & Cooling



- Gain radius defined as location where heating overtakes cooling:

$$R_g : q^- = q^+$$

Comparative competition of inputs:

- Cooling determined by matter temperature & ν emissivity:

$$q_{approx}^- \propto (k_B T)^6$$

$$q_{exact}^- \propto (k_B T)^6 \cdot [Y_p \mathcal{F}_5(\eta_e) + Y_n \mathcal{F}_5(-\eta_e)]$$

$$q_{IDSA}^- = \frac{4\pi c}{(hc)^3 \rho} \int_0^\infty d\epsilon \epsilon^3 f_{FD}(\epsilon) j_{IDSA}(\epsilon)$$

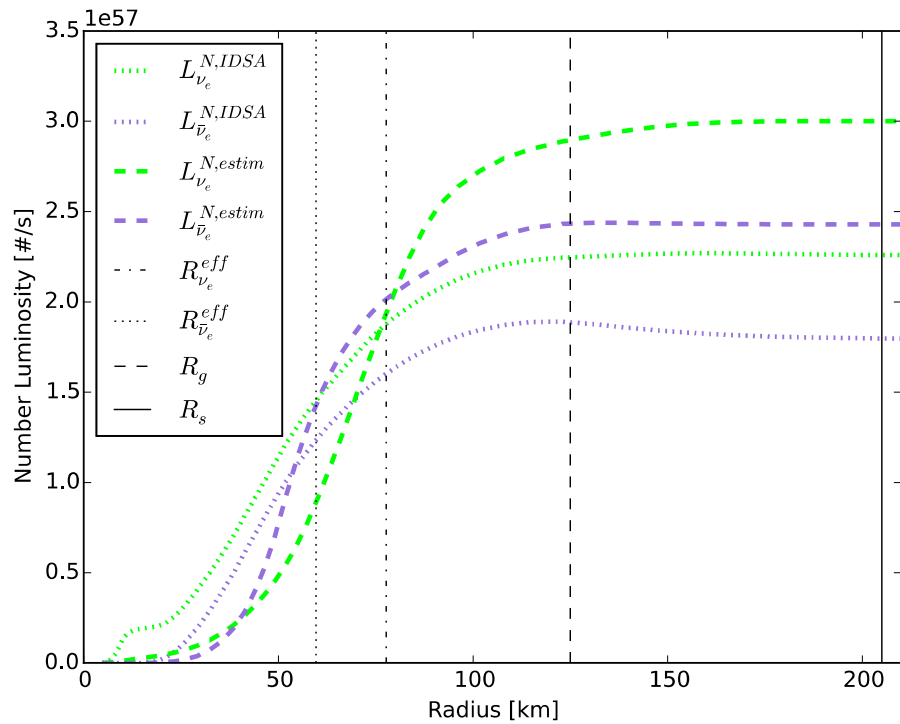
- Heating dependent on luminosity from re-absorption:

$$q_{approx}^+ \propto \frac{L_\nu}{r^2 \cdot \langle \mu_\nu \rangle} \cdot (k_B T_\nu)^2$$

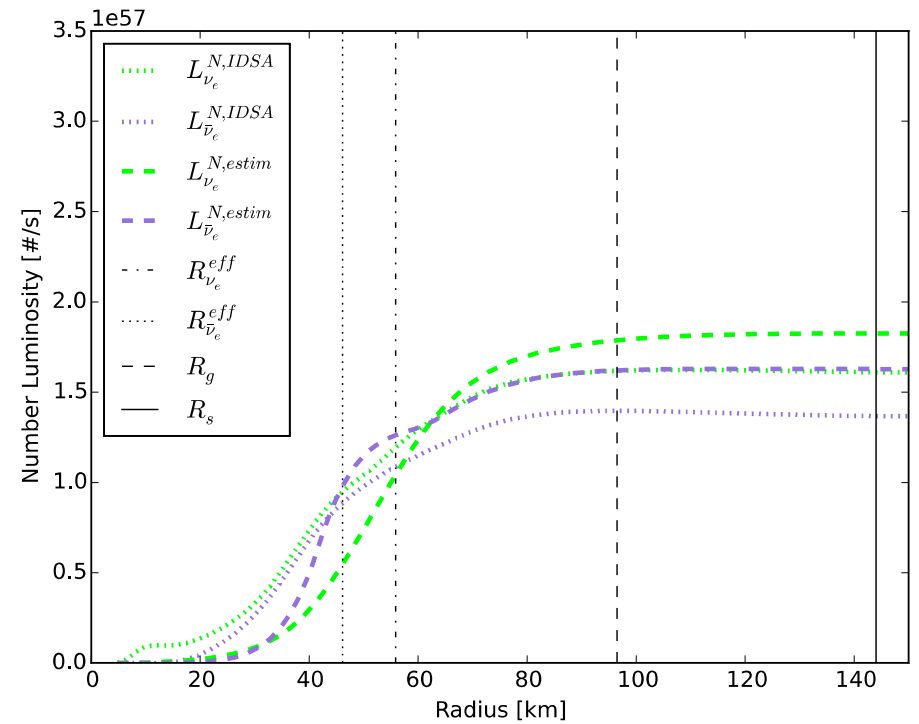
$$q_{exact}^+ \propto \frac{L_\nu}{4\pi r^2 \cdot \langle \mu_\nu \rangle} \cdot \langle \epsilon_\nu^2 \rangle \cdot (Y_n + Y_p)$$

Luminosities vs. Radii: Putting it all together...

- Number luminosity at infinity is well approximated compared to IDSA.
- Diffusion prescription spreads, or smears, L_ν^{FD} into the neutrinosphere.
- Starting point determined by Fermi blackbody light bulb!
- How does it act thereafter? $L_{\text{acc}} = L_\nu(R_g) - L_\nu(R_\nu) \quad L_{\text{acc}} \geq 0$
 $\mathcal{H} - \mathcal{C} = L_\nu(R_g) - L_\nu(R_s) \quad L_\nu(R_g) \geq L_\nu(R_s)$

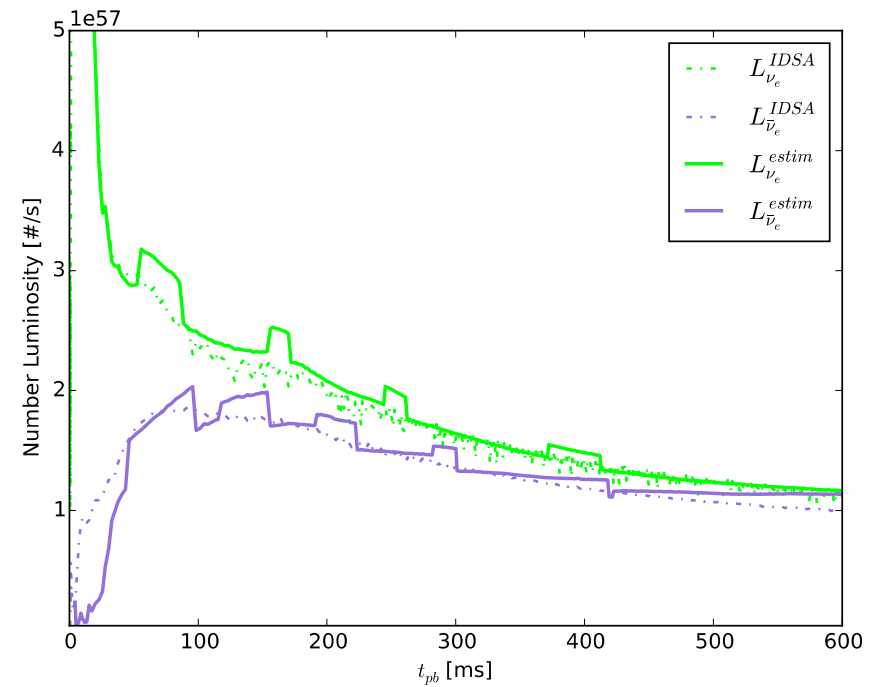
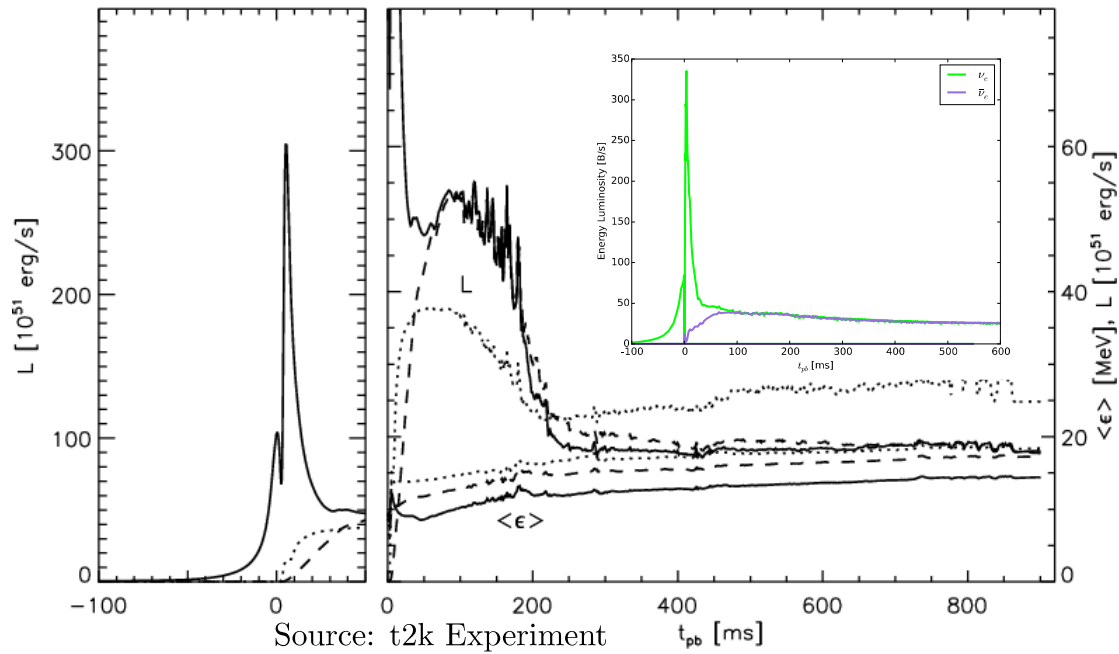


$t_{\text{pb}} \sim 120$ ms p.b.



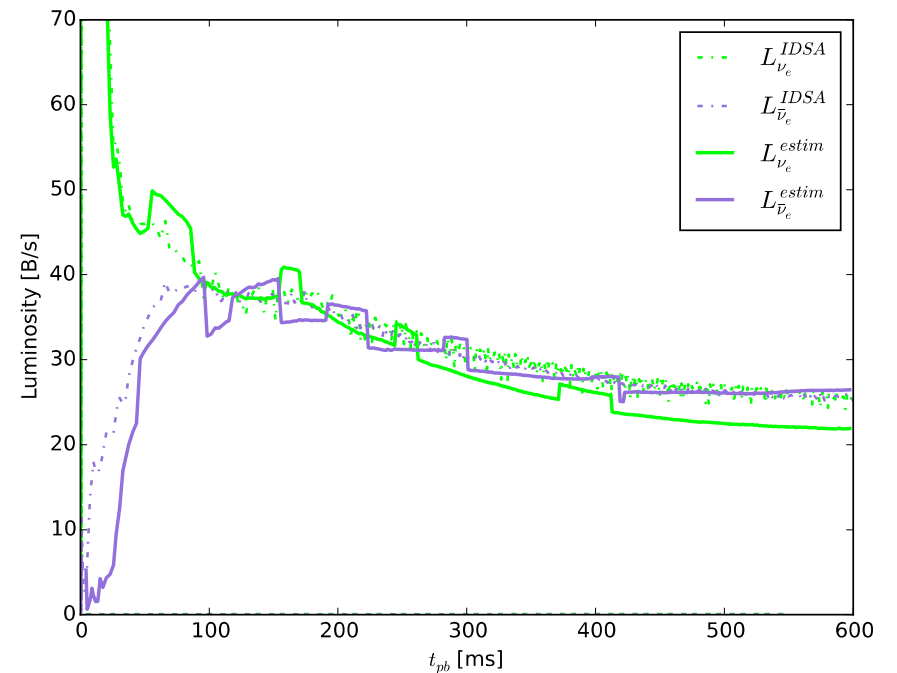
$t_{\text{pb}} \sim 300$ ms p.b.

Luminosity Estimate vs. IDSA in Time

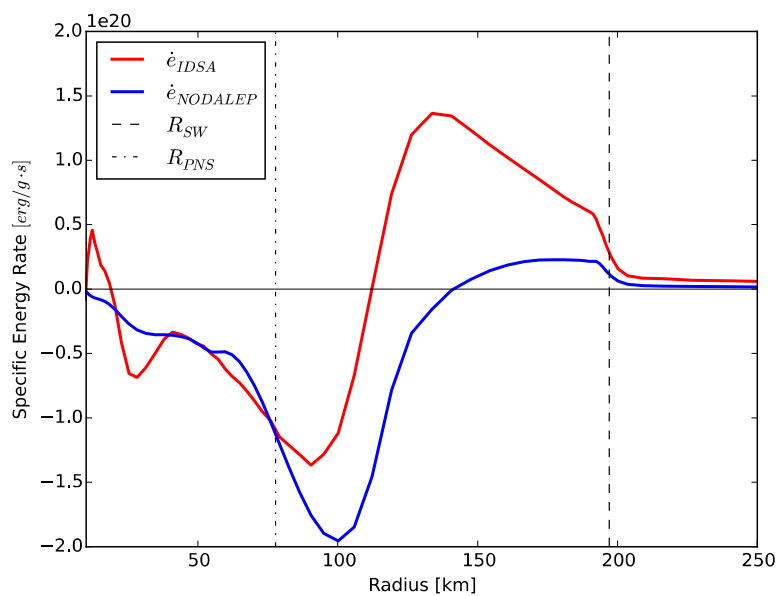
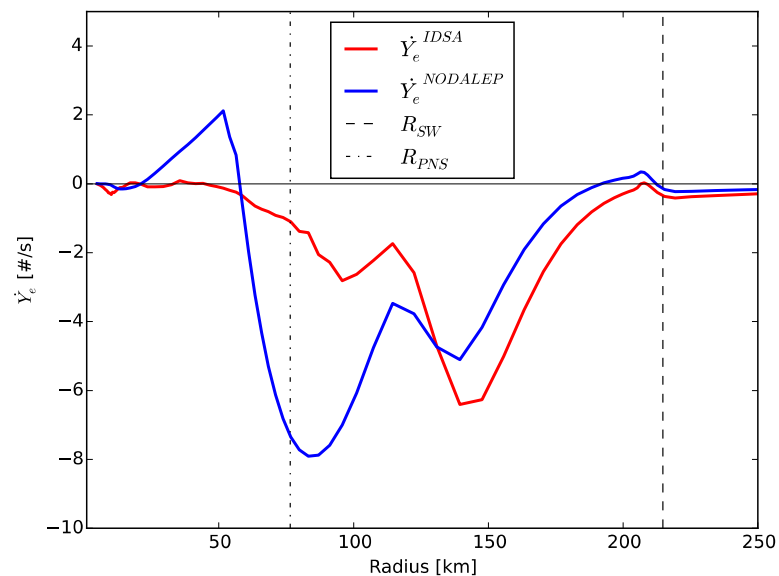
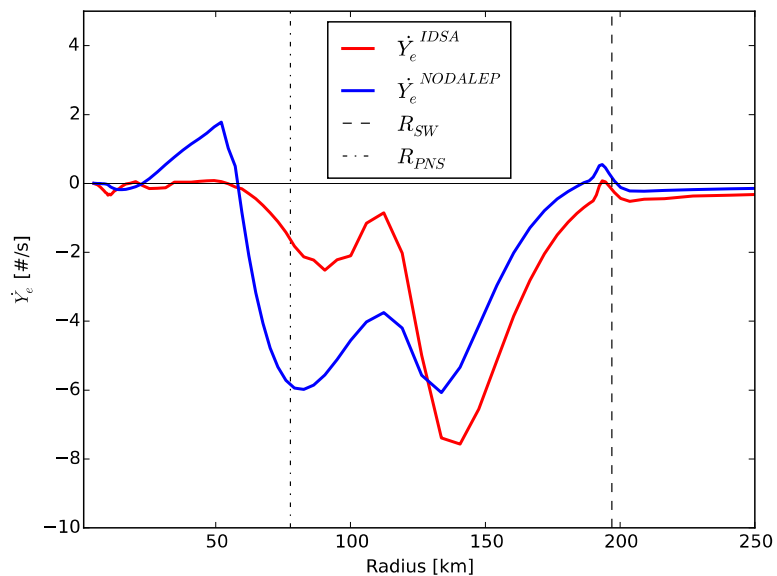


Characteristic Parts of Neutrino Signal:

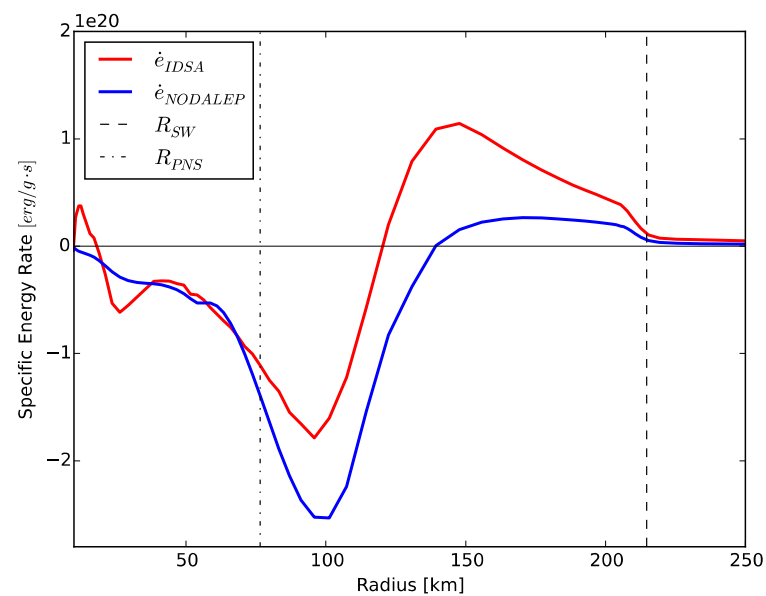
- Prompt ν_e burst
- Stalling & Luminosity from Accretion
- Late cooling of proto-neutron star



Evolution of Energy & Abundance Change Rates (\dot{e} & \dot{Y}_e): NODALEP vs. IDSA

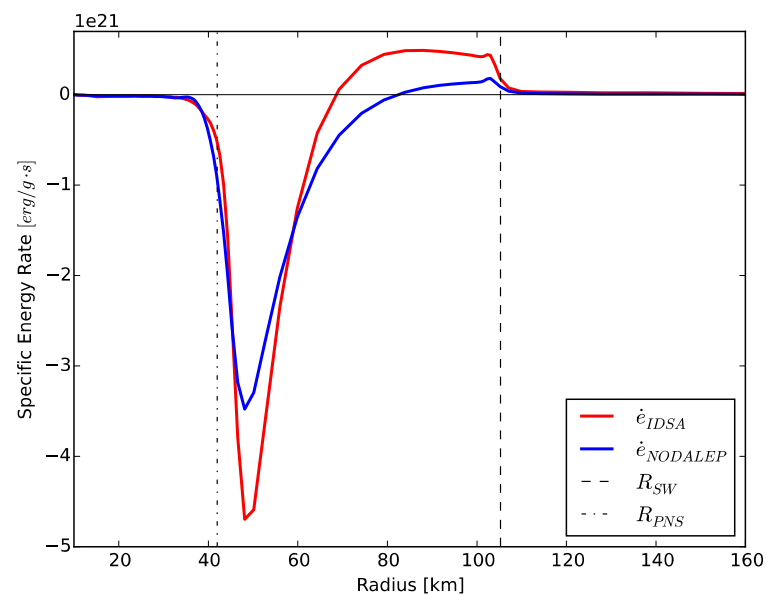
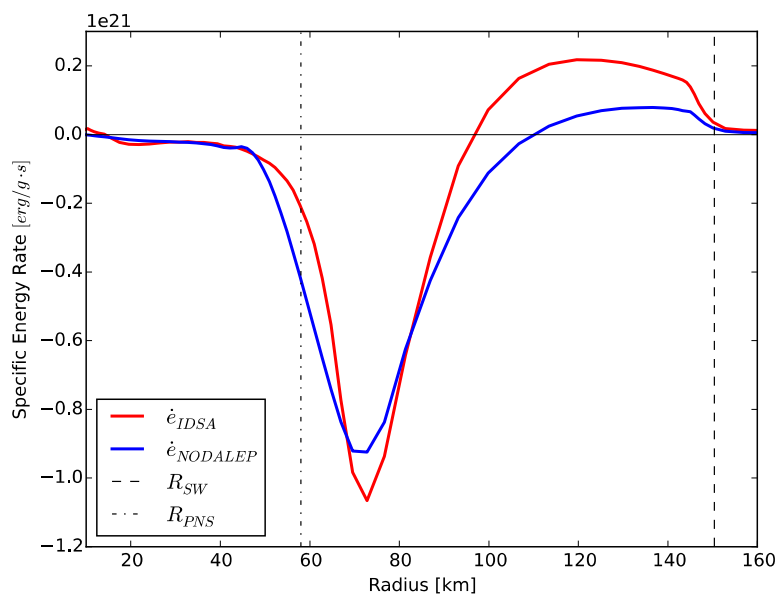
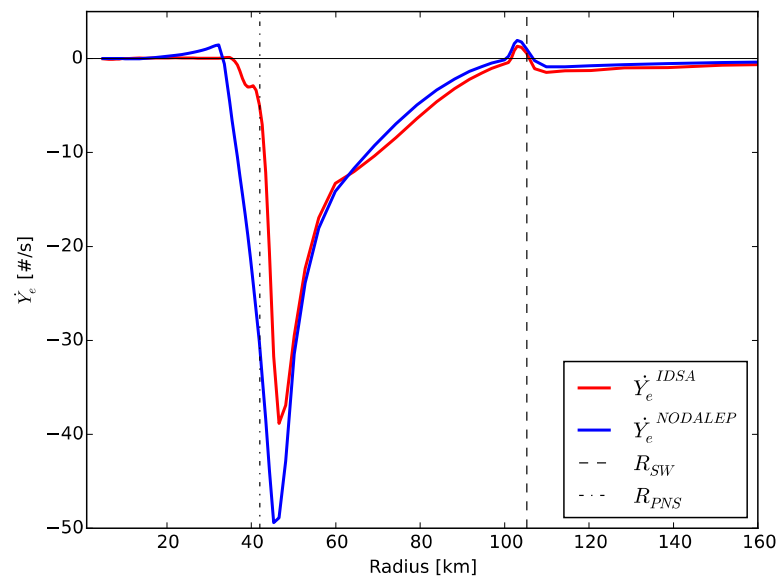
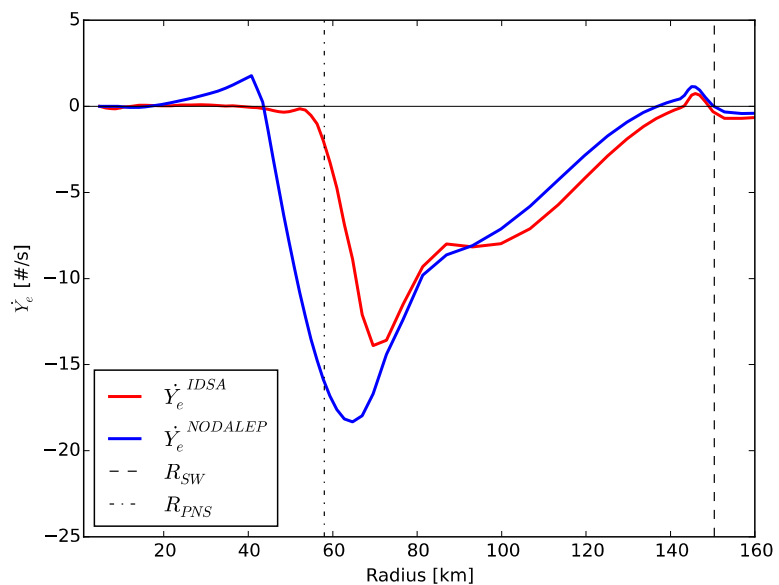


$t_{pb} \sim 100$ ms p.b.



$t_{pb} \sim 120$ ms p.b.

Evolution of Energy & Abundance Change Rates (\dot{e} & \dot{Y}_e): NODALEP vs. IDSA



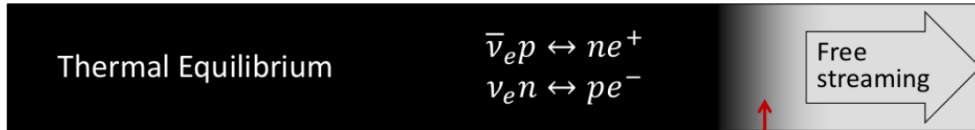
$t_{pb} \sim 300$ ms p.b.

$t_{pb} \sim 600$ ms p.b.

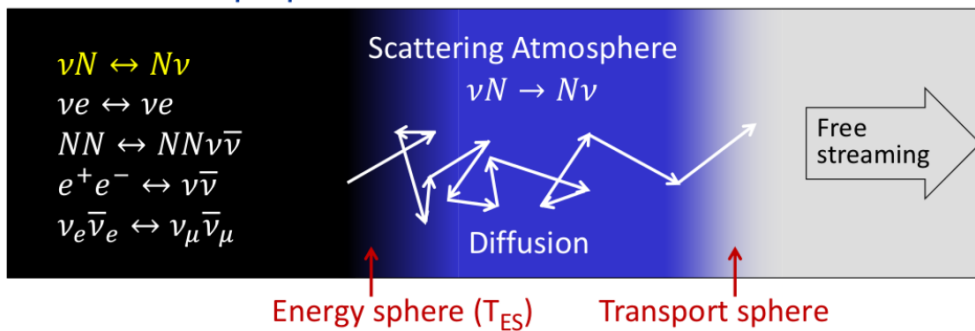
Why the difference?

- Raffelt (2012)

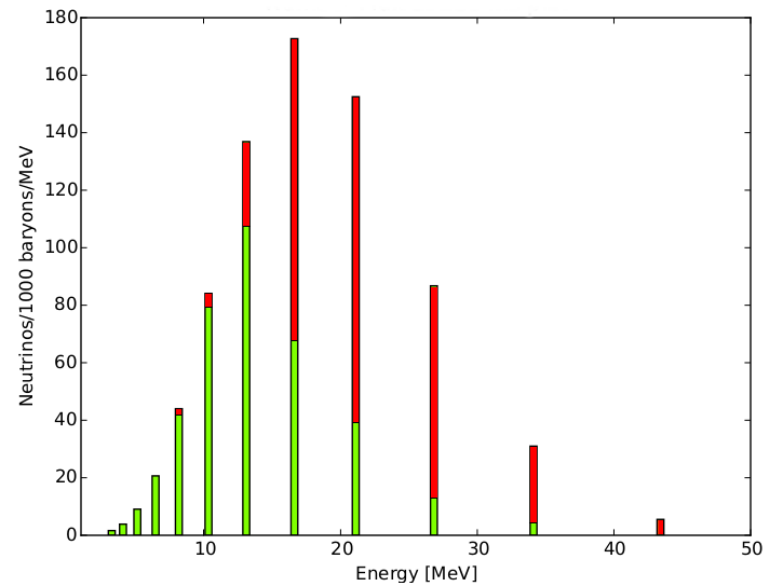
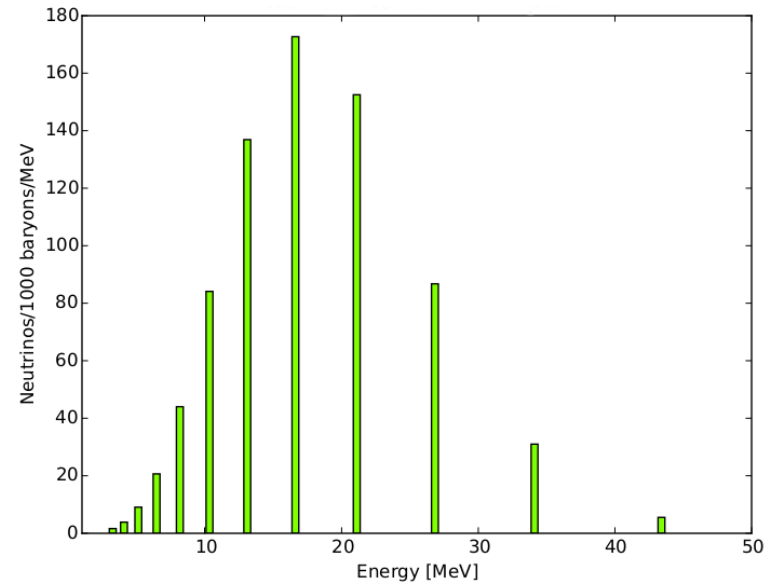
Electron flavor (ν_e and $\bar{\nu}_e$)



Other flavors ($\nu_\mu, \bar{\nu}_\mu, \nu_\tau, \bar{\nu}_\tau$)



- Spectral “clipping”

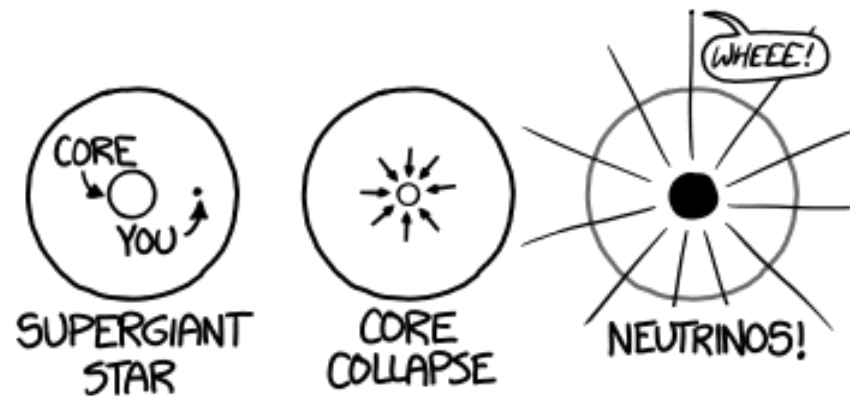


- In our new transport, we have a superposition of ν reactions, yielding a similar “diffusive atmosphere”.
- Luminosity estimates are only approximate!
- Spectral calculation would improve result.

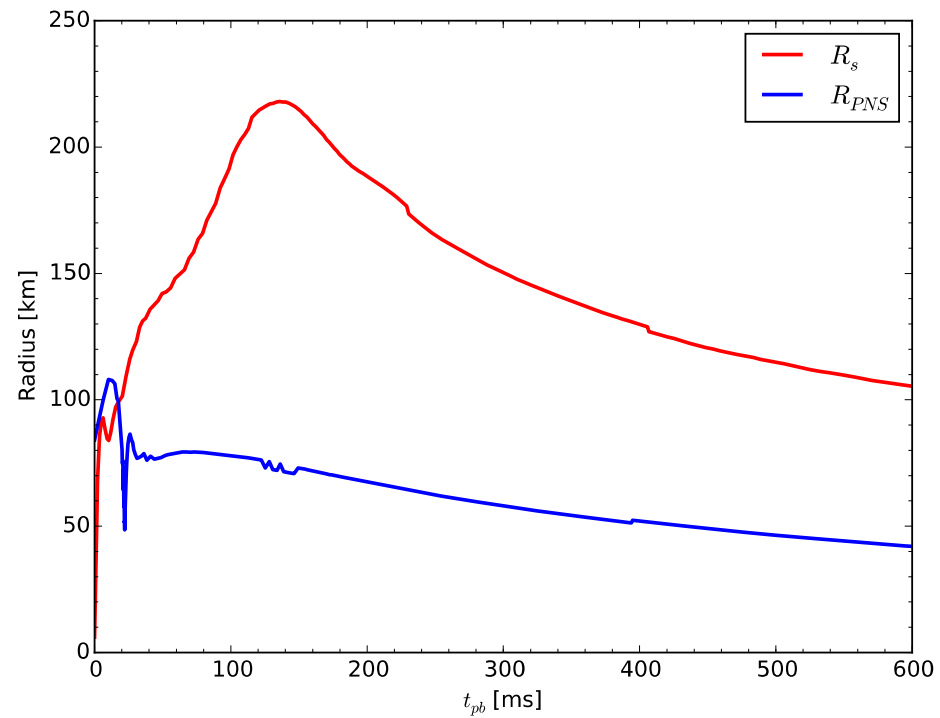
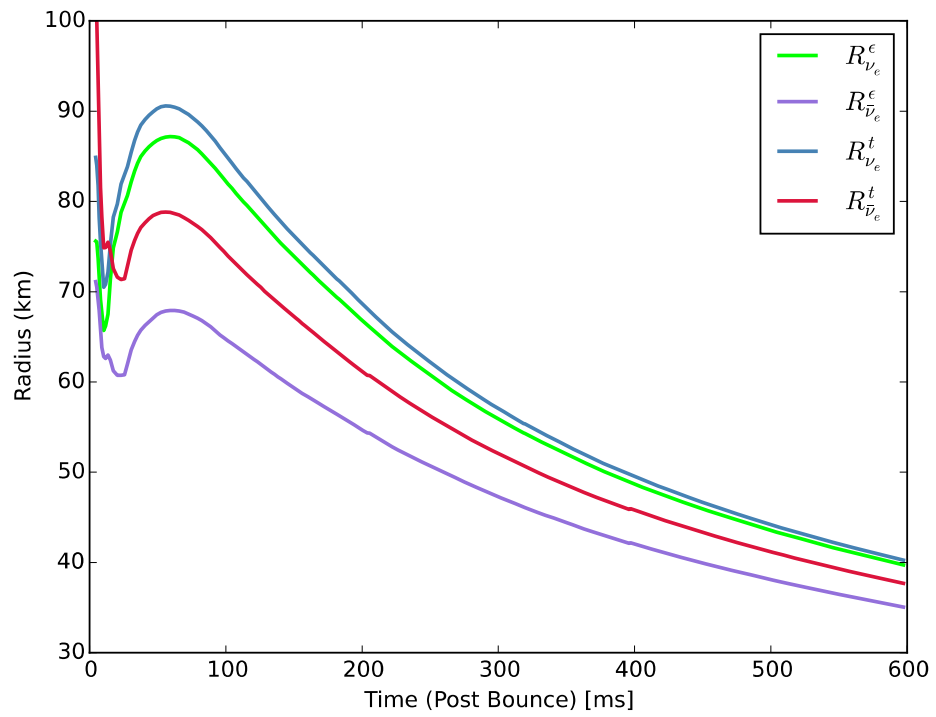
Summary & Conclusion

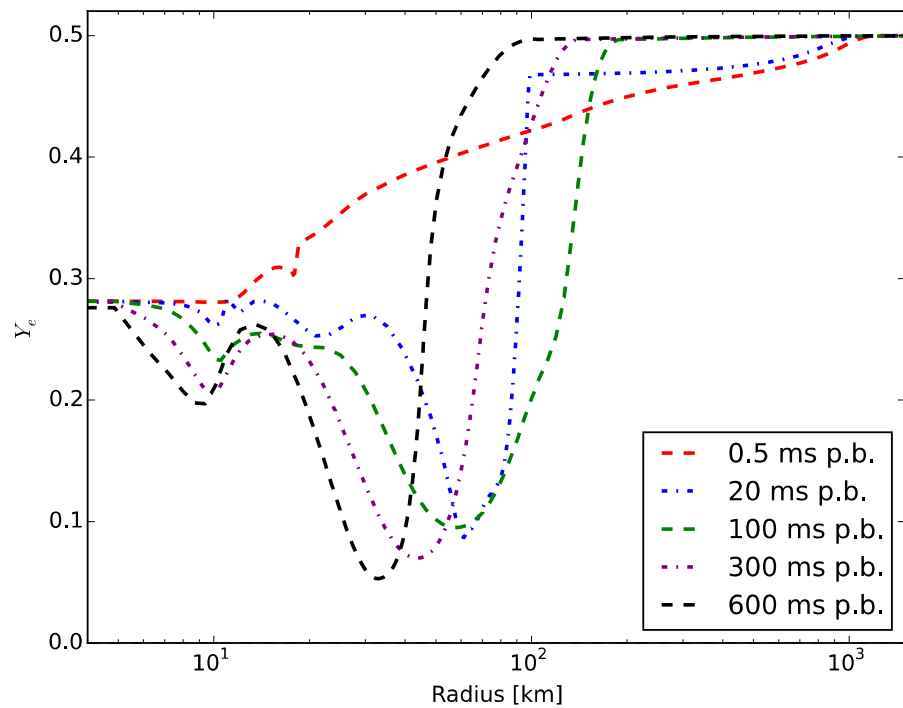
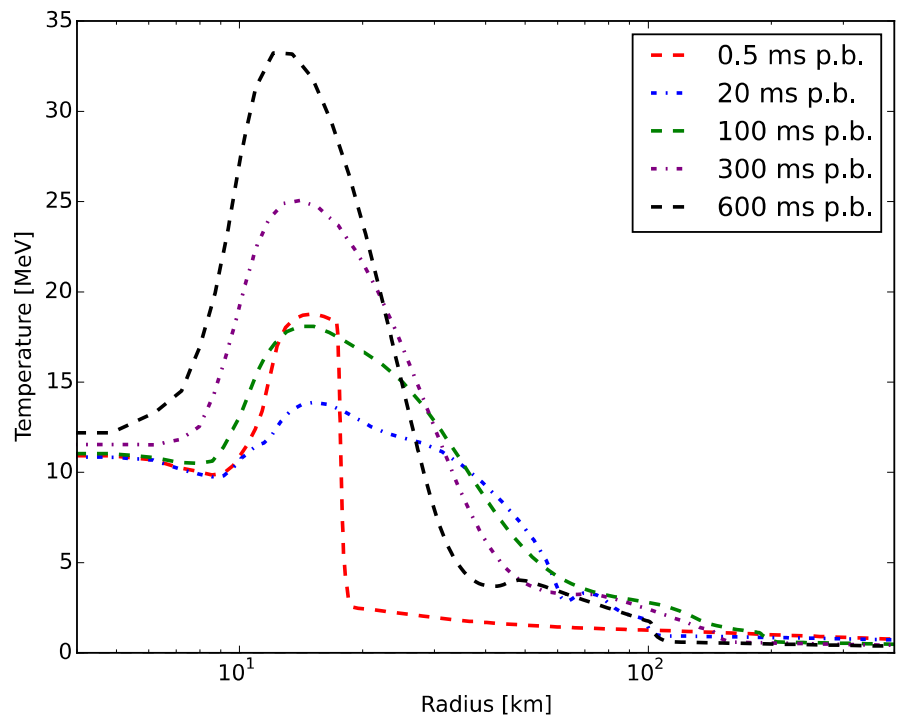
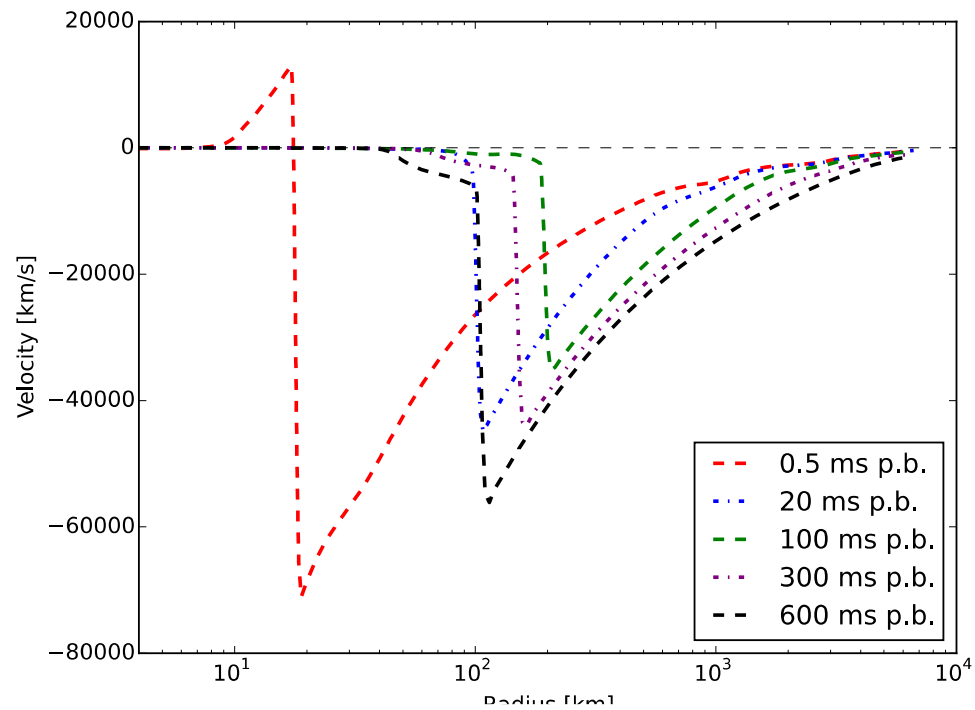
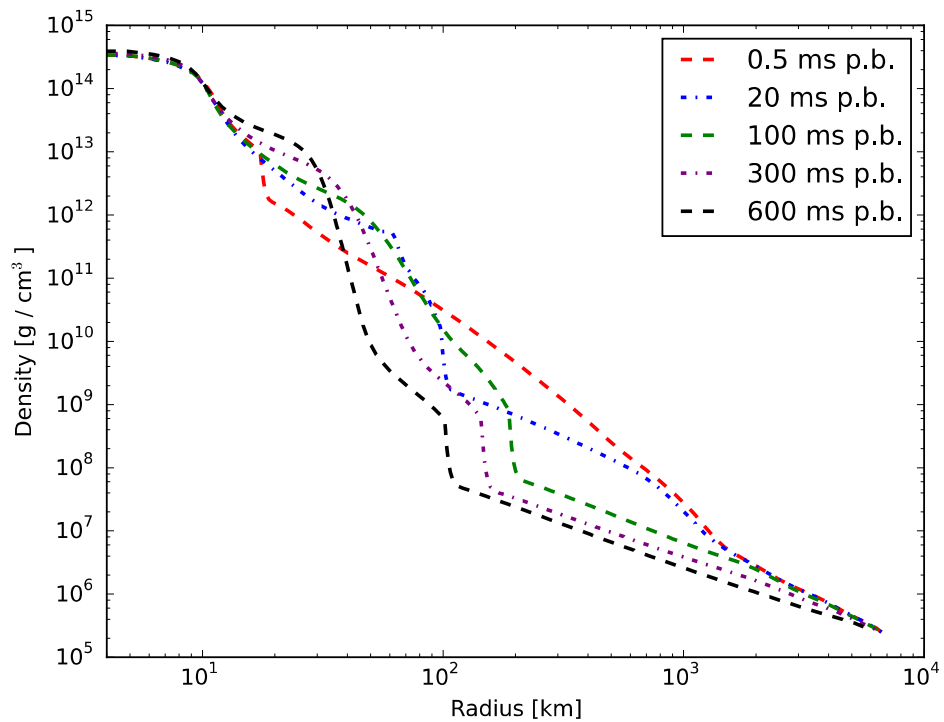
- Shock revival by neutrino heating is currently a plausible explanation for successful explosions of CCSN
- A simple, efficient neutrino treatment is demanded
- A new neutrino treatment scheme for luminosity and specific number/energy rate estimation was written and compared to output from Agile-IDSA
- Most major features of IDSA were reproduced in our neutrino treatment.
- Better treatment of neutrino degeneracy parameters at different times is needed.
- Regions between neutrino spheres contribute to greatest uncertainty
- Spectral “clipping” may motivate the need for a spectral treatment.

Thank you for your attention!



Questions?





Mean Free Path (logarithmic) [km] at 100 ms p.b.

